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New complexity analysis for primal-dual interior-point methods for self-scaled optimization problems

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Abstract

A linear optimization problem over a symmetric cone, defined in a Euclidean Jordan algebra and called a self-scaled optimization problem (SOP), is considered. We formulate an algorithm for a large-update primal-dual interior-point method (IPM) for the SOP by using a proximity function defined by a new kernel function, and we obtain the best known complexity results of the large-update IPM for the SOP by using the Euclidean Jordan algebra techniques. **MSC:** 90C51; 90C25; 65K05

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1 Introduction and preliminaries

Primal and dual interior-point methods (IPMs) have been well known as the most effective methods for solving wide classes of optimization problems, for example, the linear optimization (LO) problem, the quadratic optimization problem (QOP), the semidefinite optimization (SDO) problem, the second-order cone optimization (SOCO) problem, and the convex optimization problem (CP).

The so-called barrier update parameter θ in algorithms for IPMs plays an important role in both theory and practice of IPMs. Usually, if θ is a constant independent of the dimension of the problem, then the algorithm is called a *large-update* method. If it depends on the dimension, then the algorithm is said to be a *small-update* method. Large-update methods are much more efficient than small-update methods in practice [1], but have a worst-case iteration bound. Such a gap between theory and practice has been referred to as irony of IPMs [2]. Recently, many authors have tried to reduce the gap of the worst-case iteration bound between the large-update IPM and the small-update IPM.

Using self-regular proximity functions instead of a classical logarithmic barrier function, Peng *et al.* [3–5] improved the complexity of large-update IPMs for the LO problem, the SDO problem, and the SOCO problem. Bai *et al.* [6] introduced a new class of eligible kernel functions. The class was defined by some simple conditions on the kernel function and its derivatives. The best iteration bound for the LO problem, which was given by Bai *et al.* [6], is $\mathcal{O}(\sqrt{n} \log n \log \frac{n}{\epsilon})$. Recently, Wang *et al.* [7] obtained the complexity result $\mathcal{O}(n \log n/\epsilon)$ for the SDO problem based on a simple kernel function. Bai and Wang [8]



© 2012 Choi and Lee; licensee Springer. This is an Open Access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/2.0), which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited. obtained the best known complexity result for the SOCO problem based on a parametric kernel function including the classical logarithmic function, the prototype regular kernel function, and the non-self-regular kernel function. Very recently, using the kernel function $\phi(t) = (t^2 - 1)/2 + (e^{t^{-q}} - 1)/q$, Choi and Lee [9, 10] have obtained the complexity results of large-update primal-dual IPMs for SDO and SOCO, $\mathcal{O}(\sqrt{n}(\log n)^{(q+1)/q} \log n/\epsilon)$ and $\mathcal{O}(\sqrt{N}(\log N)^{(q+1)/q} \log N/\epsilon)$, respectively.

In this paper, we consider a linear optimization problem over a symmetric cone which is defined in a Euclidean Jordan algebra. Nesterov and Todd [11] proposed first this kind of an optimization problem under the name of convex programming for self-scaled cones and established the polynomial complexity of the primal-dual interior point method using the so-called NT (Nesterov-Todd) direction [12]. We call the linear optimization problem over the symmetric cone the self-scaled optimization problem (SOP).

Faybusovich first studied the SOP in view of a Euclidean Jordan algebra and gave a theoretical background for nondegeneracy assumptions and the uniqueness of solutions for Newton systems in IPMs for the SOP [13], presented a short-step path-following algorithm for a quadratic programming problem defined on the intersection of a symmetric cone with an affine subspace [14] and obtained complexity estimates for a long-step primal-dual interior-point algorithm for the optimization problem of the minimization of a linear function on a feasible set obtained as the intersection of an affine subspace and a symmetric cone [15]. SOPs include linear optimization problems, semidefinite optimization problems, second-order optimization problems, and various combinations of these types of problems as special cases. Schmieta and Alizadeh [16] extended primal-dual interior point algorithms for LOs, SDOs, and SOCOs to SOPs by using logarithmic barrier functions.

Baes raised an open question in his monograph [17] as follows: The theory of self-regular functions has been created for linear programming by Jiming Peng, Cornelius Roos, and Tamás Terlaky [5]. They subsequently extended it to second-order programming and semidefinite programming separately using implicitly the aforementioned construction. However, the unified treatment of this theory using the Jordan algebraic framework is not accomplished yet.

Choi and Lee [18] gave primal-dual interior point algorithms by using a very simple self-regular function $\psi(t) = \frac{1}{2}(t - \frac{1}{t})^2$, t > 0 for the SOP and gave partial answers for the question of Baes. Very recently, Vieira [19, 20] gave complete answers for the open question of Baes by proving the *e*-convexity property of eligible kernel functions and, in particular, he presented the iteration complexity results for ten eligible kernel functions. Among ten kernel functions in [19], the best iteration complexity for a large-update method was obtained for $\psi(t) = \frac{t^2-1}{2} + \frac{t^{1-q}-1}{q-1}$ with $q = \log r$, and its iteration complexity is $\mathcal{O}(\sqrt{r}\log r\log \frac{r}{\epsilon})$, which is the best known one.

In this paper, we define a new eligible kernel function $\psi(t) = \frac{t^2-1}{2} + \frac{e^{p(t^{-q}-1)}-1}{pq}$, $p \ge 1$ and $q \ge 1$ for t > 0, which was modified from the one in [9, 10], and obtain the best known iteration complexity result for the large-update IPM of the SOP by using the analysis emphasized on the kernel function and the Euclidean Jordan algebra techniques. In our algorithm, we use the well-known lemma for the upper bound of the μ -update (see Lemma 3.1) instead of using Theorem 5.4 in [20]. The lemma makes our analysis in the outer while loop easy. We refer to Theorem 4.9 and Proposition 5.6 in [20] for complexity analysis.

But we use Proposition 3.1 in [18] obtained from the technique of Sun and Sun [21] instead of using Proposition 5.7 in [20].

This paper is organized as follows. In Section 2, we introduce our kernel functions, formulate the Newton system for the SOP, and present a useful inequality for our proximity function. In Section 3, we give an algorithm for the SOP and calculate an upper bound for the proximity function after μ -update. We calculate an upper bound for difference between proximity functions after one step in inner iterations and then determine our default step size for search directions. We present a worst-case iteration bound for our large-update primal-dual interior point method for the SOP.

Now, we give definitions and preliminary properties for a Euclidean Jordan algebra which are found in [22] and will be used in the next sections.

Definition 1.1 ([22]) A finite-dimensional real vector space *V* is called an algebra if a bilinear mapping $(x, y) \rightarrow x \circ y$ from $V \times V$ to *V* is defined.

- An algebra V is called a Jordan algebra if the following hold:
- (i) commutativity: for all $x, y \in V$, $x \circ y = y \circ x$;
- (ii) Jordan's axiom: for all $x, y \in V$, $x^2 \circ (x \circ y) = x \circ (x^2 \circ y)$, where $x^2 = x \circ x$.

A Jordan algebra V is said to be Euclidean if

(iii) $x^2 + y^2 = 0 \Rightarrow x = y = 0$, equivalently, there exists an inner product $(\cdot|\cdot)$ on *V* such that $(x \circ y|z) = (y|x \circ z)$.

A Jordan algebra V is simple if it does not contain any non-trivial ideal. The Jordan algebra may not be associative, but it is power-associative, *i.e.*, $x^p \circ x^q = x^{p+q}$. We assume a Jordan algebra V has an identity element, *i.e.*, there exists e such that $x \circ e = e \circ x = x$. Since V is finite-dimensional, given $x \in V$, there exists a minimal positive integer k such that the vectors e, x, \ldots, x^k are linearly dependent. Denote this integer m(x). We define the rank of V as

 $\operatorname{rank}(V) = r = \max\{m(x) \mid x \in V\}.$

An element $x \in V$ is said to be invertible if there exists an element $y \in \mathbb{R}[x]$ such that $x \circ y = e$, where $\mathbb{R}[x]$ is the algebra over \mathbb{R} of polynomials in one variable with coefficients in \mathbb{R} . It is defined by x^{-1} . An element $v \in V$ is called idempotent if $v^2 = v$. For an element $x \in V$, let L(x) be a linear map of V defined as $L(x)y = x \circ y$. The cone of squares

$$\overline{\Omega} := \left\{ x^2 \mid x \in V \right\}$$

is a symmetric cone; the following conditions hold:

- (i) for every pair of x, y ∈ int Ω, there is an invertible linear transformation L: V → V such that L(Ω) = Ω and L(x) = y;
- (ii) $\overline{\Omega} = \overline{\Omega}$, where $\overline{\Omega} := \{y \in V \mid (x, y) \ge 0, \text{ for any } x \in \overline{\Omega}\}.$

Let $\Omega = \operatorname{int} \overline{\Omega}$. Then $\Omega = \{x^2 \mid x \in V \text{ is invertible}\} = \{x \in V \mid L(x) \text{ is positive definite}\}.$

Definition 1.2 ([22]) Let $c_1, \ldots, c_k \in V$. Then $\{c_1, \ldots, c_k\}$ is said to be a Jordan frame if c_i , $i = 1, \ldots, k$ are non-zero and cannot be written as a sum of other two idempotents, and the

following properties hold:

$$\begin{cases} c_i^2 = c_i, \\ c_i \circ c_j = 0 & \text{if } i \neq j, \\ \sum_{i=1}^k c_i = e. \end{cases}$$

Theorem 1.1 (Theorem III.1.2 in [22]) For every $x \in V$, there exist a Jordan frame $\{c_1(x), \ldots, c_r(x)\}$ and real numbers $\lambda_1(x), \ldots, \lambda_r(x)$ such that

$$x = \lambda_1(x)c_1(x) + \dots + \lambda_r(x)c_r(x).$$
(1)

The numbers $\lambda_i(x)$, for all i = 1, ..., r, are said to be the eigenvalues of x, and (1) is called the eigenvalue (or spectral) decomposition of x. Now, it is possible to extend the definition of any real-valued function $\psi(\cdot)$ to elements of the Euclidean Jordan algebra via their eigenvalues:

$$\psi(x) := \psi(\lambda_1(x))c_1(x) + \dots + \psi(\lambda_r(x))c_r(x).$$
⁽²⁾

Particularly, we have some examples as follows:

- (i) Square root: $x^{1/2} = \lambda_1^{1/2}(x)c_1(x) + \cdots + \lambda_r^{1/2}(x)c_r(x)$ if all $\lambda_i(x) \ge 0$.
- (ii) Inverse: $x^{-1} = \lambda_1^{-1}(x)c_1(x) + \cdots + \lambda_r^{-1}(x)c_r(x)$ if all $\lambda_i(x) \neq 0$.
- (iii) Square: $x^2 = \lambda_1^2(x)c_1(x) + \dots + \lambda_r^2(x)c_r(x)$.

From the above examples, we know that for $x \in \overline{\Omega}$, $\lambda_i(x^{1/2}) = \lambda_i^{1/2}(x)$ and for $x \in \Omega$, $\lambda_i(x^{-1}) = \lambda_i^{-1}(x)$. Let us denote by $\psi'(x)$ the derivative of $\psi(x)$ with respect to $\lambda_i(x)$:

$$\psi'(x) := \psi'(\lambda_1(x))c_1(x) + \dots + \psi'(\lambda_r(x))c_r(x).$$
(3)

In the Jordan algebra, we define the determinant of *x* and the trace of *x* as follows:

$$\det(x) = \prod_{i=1}^r \lambda_i(x), \qquad \operatorname{tr}(x) = \sum_{i=1}^r \lambda_i(x).$$

Since *V* is a Euclidean Jordan algebra, $\langle x, y \rangle := tr(x \circ y)$ is a scalar product on *V* (see Proposition III.1.5 in [22]). The following lemma is called the second Pierce decomposition theorem which will be used in Section 3.

Lemma 1.1 (Theorem IV.2.1 in [22], Theorem 2.6.6 (Second Pierce decomposition theorem) in [17]) *Let* $\{c_1, \ldots, c_r\}$ *be a Jordan frame of V*. *If*

$$V_{ij} := \begin{cases} \{v_{ij} \mid c_i \circ v_{ij} = v_{ij}\} & \text{if } i = j, \\ \{v_{ij} \mid c_i \circ v_{ij} = \frac{1}{2}v_{ij}\} \cap \{v_{ij} \mid c_j \circ v_{ij} = \frac{1}{2}v_{ij}\} & \text{if } i \neq j, \end{cases}$$

we have

- (i) $V = \bigoplus_{1 \le i \le j \le r} V_{ij};$
- (ii) $V_{ij} \circ V_{kl} = \{0\}, if\{i, j\} \cap \{k, l\} = \emptyset;$
- (iii) $V_{ij} \circ V_{jk} \subset V_{ik}$, if $i \neq k$;
- (iv) $\operatorname{tr}(v_{ik}) = 0$, for $v_{ik} \in V_{ik}$ if $i \neq k$.

Consider the following self-scaled optimization problem (SOP):

(P) Minimize
$$\langle c, x \rangle$$

subject to $\langle a_i, x \rangle = b_i, \quad i = 1, ..., m,$
 $x \in \overline{\Omega},$

and its dual problem:

(D) Maximize
$$\sum_{\substack{i=1\\m}}^{m} b_i y_i$$

subject to
$$\sum_{\substack{i=1\\s\in\overline{\Omega}, \quad y\in\mathbb{R}^m}}^{m} y_i a_i + s = c,$$

where $c, a_1, \ldots, a_m \in V$ and $b \in \mathbb{R}^m$ are given. We call $x \in \overline{\Omega}$ primal feasible if $\langle a_i, x \rangle = b_i$ for $i = 1, \ldots, m$. Similarly, $(y, s) \in \mathbb{R}^m \times \overline{\Omega}$ is called dual feasible if $\sum_{i=1}^m y_i a_i + s = c$. Let $Ax = (\langle a_1, x \rangle, \ldots, \langle a_m, x \rangle)^T$ for any $x \in V$. Then $A : V \to \mathbb{R}^m$ is a linear transformation. Throughout this paper, we assume that A is surjective. Then its adjoint A^T is injective and $A^T y = \sum_{i=1}^m y_i a_i$, where $y = (y_1, \ldots, y_m)^T \in \mathbb{R}^m$. So, we can reformulate (P) and (D) as follows:

(P) Minimize $\langle c, x \rangle$ subject to Ax = b, $x \in \overline{\Omega}$,

and its dual problem:

(D) Maximize
$$b^T y$$

subject to $A^T y + s = c$,
 $s \in \overline{\Omega}, \quad y \in \mathbb{R}^m$

We can check that weak duality between (P) and (D) holds, that is, $inf(P) \ge sup(D)$. From now on, we assume that both (P) and (D) satisfy the interior-point condition (IPC), that is, there exists (x^0, y^0, s^0) such that $Ax^0 = b$, $x^0 \in \Omega$, $A^Ty^0 + s^0 = c$, $s^0 \in \Omega$. Then there exists a pair of optimal solutions (x, y, s) of (P) and (D), and inf(P) = sup(D) [11, 23].

The following lemma is well known [13, 17, 22, 24].

Lemma 1.2 For $x, s \in V$, the following statements are equivalent:

(i) $x, s \in \overline{\Omega} \text{ and } \langle x, s \rangle = 0;$ (ii) $x, s \in \overline{\Omega} \text{ and } x \circ s = 0.$

Using Lemma 1.2, we can check (see Proposition 2.1 in [13]) that finding a pair of optimal solutions (x, y, s) of (P) and (D) is equivalent to solving the following Newton system:

$$\begin{aligned} Ax &= b, \\ A^T y + s &= c, \\ x \circ s &= 0, \\ x, s &\in \overline{\Omega}, \qquad y \in \mathbb{R}^m. \end{aligned} \tag{4}$$

The basic idea of primal-dual IPMs is to replace the third equation in (4), the so-called *complementarity condition* for the SOP, by the parameterized system with a positive parameter μ :

$$Ax = b,$$

$$A^{T}y + s = c,$$

$$x \circ s = \mu e,$$

$$x, s \in \Omega, \qquad y \in \mathbb{R}^{m}.$$
(5)

For each $x \in V$, we define the quadratic representation as follows:

$$Q_x := 2L^2(x) - L(x^2).$$

Lemma 1.3 ([16]) Let $x, s \in \Omega$ and p be invertible. Then $x \circ s = \mu e$ if and only if $Q_p x \circ Q_{p^{-1}} s = \mu e$.

Proposition 1.1 (Proposition 18 in [16]) *If* $x, s \in \Omega$, *then* $Q_x s \in \Omega$.

Let $x, s \in \Omega$. Then there uniquely exists $p \in \Omega$ such that $Q_{p^2}x = s$ [25, 26]. So, we can choose $p \in \Omega$ such that $Q_px = Q_{p^{-1}s}$. Such a choice exists and is unique, and leads to the Nesterov-Todd (NT) method.

From Lemma 1.3, the system (5) becomes

$$\begin{cases}
Ax = b, \\
A^{T}y + s = c, \\
Q_{p}x \circ Q_{p^{-1}}s = \mu e, \\
x, s \in \Omega, \quad y \in \mathbb{R}^{m}.
\end{cases}$$
(6)

Then, for each $\mu > 0$, the parameterized system (6) has a unique solution $(x(\mu), y(\mu), s(\mu))$ [11, 27], which is called a μ -center of (P) and (D). The set of μ -centers, that is, $C = \{(x(\mu), y(\mu), s(\mu)) \mid \mu > 0\}$, is said to be the *central path* of (P) and (D). Therefore, as μ tends to zero, $(x(\mu), y(\mu), s(\mu))$ converges to a pair of optimal solutions of (P) and (D) [13, 28].

In general, IPMs for the SOP consist of two strategies. The first one, which is called the inner iteration scheme, is to keep the iterative sequence in a certain neighborhood of the central path or to keep the iterative sequence in a certain neighborhood of the μ -center. And the second one, called the outer iteration scheme, is to decrease the parameter μ to $\mu_+ := (1 - \theta)\mu$ for some $\theta \in (0, 1)$.

2 Proximity functions and search directions

Newton's method is a well-known procedure to solve a system of nonlinear equations. Most IPMs for solving the SOP employ different search directions together with suitable strategies for following the central path appropriately.

Assume that a starting point (x^0, s^0) in a certain neighborhood of the central path corresponding to $\mu = 1$ is available. We then decrease μ to $\mu_+ := (1 - \theta)\mu$ for some fixed $\theta \in (0, 1)$

and linearize the Newton system for (6) by replacing *x*, *y*, *s* with $x_+ := x + \Delta x$, $y_+ := y + \Delta y$, $s_+ := s + \Delta s$, respectively. Then we get the following system in [16]:

$$\begin{cases}
A\Delta x = 0, \\
A^{T}\Delta y + \Delta s = 0, \\
Q_{p}x \circ Q_{p^{-1}}\Delta s + Q_{p}\Delta x \circ Q_{p^{-1}}s = \mu_{+}e - Q_{p}x \circ Q_{p^{-1}}s.
\end{cases}$$
(7)

To describe our new search direction, we need more notations:

$$\bar{A} := \frac{1}{\sqrt{\mu}} A Q_{p^{-1}}, \qquad \nu := \frac{1}{\sqrt{\mu}} Q_p x = \frac{1}{\sqrt{\mu}} Q_{p^{-1}} s,
dx := \frac{1}{\sqrt{\mu}} Q_p \Delta x, \qquad ds := \frac{1}{\sqrt{\mu}} Q_{p^{-1}} \Delta s.$$
(8)

In this case,

$$p = \left[Q_{x^{1/2}}(Q_{x^{1/2}}s)^{-1/2}\right]^{-1/2} = \left[Q_{s^{-1/2}}(Q_{s^{1/2}}x)^{1/2}\right]^{-1/2}.$$
(9)

From Proposition 1.1, $\nu \in \Omega$. Hence, $L(\nu)$ is positive definite. Thus, the system (7) is equivalent to the following system:

$$\begin{cases} \bar{A} dx = 0, \\ \bar{A}^T \Delta y + ds = 0, \\ dx + ds = v^{-1} - v. \end{cases}$$
(10)

We say that the above $(dx, \Delta y, ds)$ is called the NT search direction for the SOP. Furthermore, $\langle dx, ds \rangle = 0$, which is coming from the first and second equations of (10) or from the orthogonality of Δx and Δs .

For our IPM, we use the following new eligible kernel function:

$$\psi(t) = \frac{t^2 - 1}{2} + \frac{e^{p(t^{-q} - 1)} - 1}{pq}, \quad p \ge 1 \text{ and } q \ge 1 \text{ for } t > 0.$$
(11)

Please see the definition of an eligible function in [6]. The new kernel function (11) satisfies

$$\psi''(t) > 1$$
, $\psi'''(t) < 0$ and $\lim_{t \to 0^+} \psi(t) = \lim_{t \to \infty} \psi(t) = \infty$.

Note that $\psi(1) = \psi'(1) = 0$. Then $\psi(t)$ is determined:

$$\psi(t) = \int_{1}^{t} \int_{1}^{\xi} \psi''(\zeta) \, d\zeta \, d\xi.$$
(12)

The proximity function (measure) for (P) and (D) is

$$\Phi(x,s;\mu) := \Psi(\nu) := \operatorname{tr}(\psi(\nu)) = \sum_{i=1}^{r} \psi(\lambda_i(\nu)),$$
(13)

where $\psi(v)$ is defined by (2). Note that $\Psi(v) = 0$, if v = e (*i.e.*, $x \circ s = \mu e$) and $\Psi(v) > 0$, otherwise. Replacing the right-hand side of the last equation in (10) by $-\psi'(v)$, we have the following system from (10):

$$\begin{cases} \bar{A} dx = 0, \\ \bar{A}^T \Delta y + ds = 0, \\ dx + ds = -\psi'(v). \end{cases}$$
(14)

Let $X = \{x \in V \mid \overline{A}x = 0\}$. Then $X^{\perp} = \{\overline{A}^T y \mid y \in \mathbb{R}^m\}$. Hence, the system (14) has a unique solution. We introduce the norm-based proximity measure as follows:

$$\sigma := \|dx + ds\| = \|\psi'(v)\| = \sqrt{\|dx\|^2 + \|ds\|^2}.$$
(15)

The following lemma gives a lower bound of σ in terms of $\Psi(\nu)$.

Lemma 2.1 *For any* $v \in \Omega$ *,*

$$\sigma \geq \sqrt{2\Psi(\nu)}.$$

Proof Since (11) satisfies $2\psi(t) \leq (\psi'(t))^2$ and $\sigma^2 = \sum_{i=1}^r (\psi'(\lambda_i(v)))^2$,

$$2\Psi(\nu) \leq \sigma^2$$
.

This completes the proof.

Also, our new kernel function (11) satisfies the following exponential convexity property.

Lemma 2.2 *Let* $t_1 > 0$ *and* $t_2 > 0$ *. Then*

$$\psi(\sqrt{t_1t_2}) \leq rac{1}{2} (\psi(t_1) + \psi(t_2)).$$

The following proposition can be found in [20], but for the completeness, we give its proof.

Proposition 2.1 (Theorem 4.9 in [20]) Let Ψ be the proximity function defined in (13), then for any $x, s \in \Omega$,

$$\Psi((Q_{x^{1/2}}s)^{1/2}) \leq \frac{1}{2}(\Psi(x) + \Psi(s)).$$

Proof Since $Q_{x^{1/2}}s \in \Omega$,

$$\lambda_i ((Q_{x^{1/2}}s)^{1/2}) = \lambda_i^{1/2} (Q_{x^{1/2}}s) \text{ and } \Psi ((Q_{x^{1/2}}s)^{1/2}) = \sum_{i=1}^r \psi (\lambda_i^{1/2} (Q_{x^{1/2}}s)).$$

By Theorem 3.5 in [20],

$$\prod_{i=1}^k \lambda_i(Q_{x^{1/2}}s) \leq \prod_{i=1}^k \lambda_i(x)\lambda_i(s), \quad \text{for } k = 1, \dots, r-1,$$

and

$$\prod_{i=1}^r \lambda_i(Q_{x^{1/2}}s) = \prod_{i=1}^r \lambda_i(x)\lambda_i(s).$$

Thus,

$$\prod_{i=1}^{k} \lambda_i^{1/2}(Q_{x^{1/2}}s) \leq \prod_{i=1}^{k} \lambda_i^{1/2}(x) \lambda_i^{1/2}(s), \quad \text{for } k = 1, \dots, r-1,$$

and

$$\prod_{i=1}^r \lambda_i^{1/2}(Q_{x^{1/2}}s) = \prod_{i=1}^r \lambda_i^{1/2}(x)\lambda_i^{1/2}(s).$$

Let $\alpha_i = \lambda_i^{1/2}(Q_{x^{1/2}}s)$ and $\beta_i = \lambda_i^{1/2}(x)\lambda_i^{1/2}(s)$. Then $\alpha_i > 0$ and $\beta_i > 0$. Moreover, since these conditions satisfy the assumptions of Corollary 3.3.10 in [29] and (iii) in Corollary 3.3.10 in [29] with our kernel function (11),

$$\sum_{i=1}^{r} \psi \left(\lambda_{i}^{1/2}(Q_{x^{1/2}}s) \right) \leq \sum_{i=1}^{r} \psi \left(\lambda_{i}^{1/2}(x) \lambda_{i}^{1/2}(s) \right).$$

By Lemma 2.2, we obtain the following result:

$$\sum_{i=1}^r \psi\left(\lambda_i^{1/2}(x)\lambda_i^{1/2}(s)\right) \leq \frac{1}{2} \left(\sum_{i=1}^r \psi\left(\lambda_i(x)\right) + \sum_{i=1}^r \psi\left(\lambda_i(s)\right)\right) = \frac{1}{2} \left(\Psi(x) + \Psi(s)\right).$$

3 Algorithm and its complexity analysis

Now, we explain our algorithm for the large-update primal-dual IPM for the SOP. Assuming that a starting point in a certain neighborhood of the central path is available, we can set out from this point. Then, we will go to the outer 'while loop'. If μ satisfies $r\mu \ge \epsilon$, then it is reduced by the factor $1 - \theta$, where $\theta \in (0, 1)$. Then, we make use of the inner 'while loop', and we repeat the procedure until we find iterates that are 'close' to $(x(\mu), y(\mu), s(\mu))$, that is, the proximity $\Phi(x, s; \mu) < \tau$. Here, we apply Newton's method targeting at the new μ -centers to decide a search direction $(\Delta x, \Delta y, \Delta s)$. We return to the outer 'while loop'. The whole process is repeated until μ is small enough, say until $r\mu < \epsilon$.

The choice of the step size α is another crucial issue in the analysis of the algorithm. It has to be taken so that the closeness of the iterates to the current μ -center can improve by a sufficient amount. In the algorithm, the inner 'while loop' is called the *inner iteration* and the outer 'while loop' is called the *outer iteration*. Each outer iteration consists of an update of the parameter μ and a sequence of (one or more) inner iterations. The total number of inner iterations is the worst-case iteration bound for our algorithm.

The algorithm for our large-update primal-dual IPM for the SOP is given as follows:

Primal-dual algorithm for the SOP

Inputs

A proximity parameter $\tau > 1$; an accuracy parameter $\epsilon > 0$;

```
a variable damping factor \alpha;
   a fixed barrier update parameter \theta \in (0, 1);
   (x^0, s^0) and \mu^0 = 1 such that \Phi(x^0, s^0; \mu^0) \leq \tau.
begin
   x := x^0; s := s^0; \mu := \mu^0;
   while r\mu \geq \epsilon do
   begin
      \mu := (1 - \theta)\mu;
      while \Phi(x, s; \mu) \ge \tau do
      begin
          Solve the system (14) for \Delta x, \Delta y, \Delta s;
          Determine a step size \alpha;
         x := x + \alpha \Delta x;
         y := y + \alpha \Delta y;
         s := s + \alpha \Delta s;
      end
   end
end
```

3.1 Bound of the proximity function after μ -update

We have $\Psi(\nu) \leq \tau$ before the update of μ with the factor $1 - \theta$ at the start of each outer iteration. After updating μ in an outer iteration, the vector ν is divided by the factor $\sqrt{1-\theta}$, which in general leads to an increase in the value of $\Psi(\nu)$. Then during the inner iteration, the value of $\Psi(\nu)$ decreases until it passes the threshold τ .

As we mentioned, our kernel function (11) is eligible. To obtain an upper bound for a μ -updated proximity function in each outer iteration in the algorithm, we use the well-known Lemma 3.1, which can be induced from the decreasing part of the kernel function, instead of using theorems which can be obtained from some properties for eligible functions (for example, Theorem 3.2 in [6] and Theorem 5.4 in [20]). Both of the following lemmas make our analysis in the outer while loop easy. And we will show a theorem that an upper bound for $\Psi(\frac{1}{\sqrt{1-\theta}}\nu)$ is expressed with $\Psi(\nu)$ by using the following two lemmas.

Lemma 3.1 Let $\beta \ge 1$. Then

$$\psi(\beta t) \leq \psi(t) + \frac{(\beta^2 - 1)}{2}t^2.$$

Proof Define $\psi_b(t) := \frac{e^{p(t^{-q}-1)}-1}{pq}$. Then $\psi_b(t)$ is monotonically decreasing in *t*. So, we can easily obtain

$$\begin{split} \psi(\beta t) &= \frac{\beta^2 t^2 - 1}{2} + \psi_b(\beta t) = \frac{t^2 - 1}{2} + \psi_b(t) + \frac{\beta^2 t^2 - t^2}{2} + \psi_b(\beta t) - \psi_b(t) \\ &\leq \psi(t) + \frac{(\beta^2 - 1)}{2} t^2. \end{split}$$

Lemma 3.2 *For any* $v \in \Omega$ *, then*

$$\|\nu\|^2 \leq 2(\Psi(\nu) + 2r).$$

Proof Since $\frac{e^{p(t^{-q}-1)}}{pq}$ is positive and $pq \ge 1$, the kernel function (11) has a lower bound as follows:

$$\psi(t) \ge \frac{t^2 - 1}{2} - \frac{1}{pq} \ge \frac{t^2}{2} - 1 - 1.$$

This implies $\frac{1}{2} \sum_{i=1}^{r} \lambda_i^2(v) \leq \Psi(v) + 2r$.

Theorem 3.1 Let θ be such that $0 < \theta < 1$. Then, for any $v \in \Omega$,

$$\Psi\left(\frac{1}{\sqrt{1-\theta}}\nu\right) \leq \frac{2}{1-\theta}(\Psi(\nu)+r).$$

Proof From Lemma 3.1 with $\beta = \frac{1}{\sqrt{1-\theta}}$ and Lemma 3.2,

$$\begin{split} \Psi\bigg(\frac{1}{\sqrt{1-\theta}}\nu\bigg) &= \sum_{i=1}^r \psi\bigg(\frac{1}{\sqrt{1-\theta}}\lambda_i(\nu)\bigg) \leq \Psi(\nu) + \frac{1}{2}\bigg(\frac{1}{1-\theta} - 1\bigg)\|\nu\|^2 \\ &\leq \Psi(\nu) + \frac{\theta}{1-\theta}\big(\Psi(\nu) + 2r\big) \leq \frac{2}{1-\theta}\big(\Psi(\nu) + r\big), \end{split}$$

the last inequality comes from $\theta \in (0, 1)$.

By the assumption $\Psi(\nu) \leq \tau$ just before the update of μ ,

$$\Psi\left(\frac{1}{\sqrt{1-\theta}}\nu\right) \leq \frac{2}{1-\theta}(\tau+r).$$

We define

$$L(r,\theta,\tau)=\frac{2}{1-\theta}(\tau+r).$$

Since $\tau = \mathcal{O}(r)$ and $\theta = \Theta(1)$,

$$L = \mathcal{O}(r).$$

3.2 Determining a default step size

In this section, we compute the feasible step size α such that the proximity function is decreasing and is bound for the decrease during inner iterations; then we give our default step size $\bar{\alpha}$; $\bar{\alpha} = (3(1 + 3\sigma(1 + pq + q)(1 + p^{-1}\log 3\sigma)^{(q+1)/q}))^{-1}$. We will show that the step size not only keeps the iterates feasible but also gives rise to a sufficiently large decrease in the barrier function $\Psi(\nu)$ in each inner iteration. Let us denote the difference between the proximity before and after one step by a function of the step size, that is,

$$f(\alpha) := \Psi(\nu_+) - \Psi(\nu). \tag{16}$$

The main task in the rest of this section is to study the decreasing behavior of $f(\alpha)$.

Now, in equation (16), v_+ and p_+ are determined by x, s in (9) and (8) replaced by $x_+ := x + \alpha \Delta x$, $s_+ := s + \alpha \Delta s$, respectively, which is as follows:

$$\nu_+ := \frac{1}{\sqrt{\mu}}Q_{p_+}(x+\alpha\Delta x) = \frac{1}{\sqrt{\mu}}Q_{p_+^{-1}}(s+\alpha\Delta s).$$

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Lemma 3.3 (Proposition II.3.3 in [22]) Let x and s be elements in V. Then

- (i) $(Q_x s)^{-1} = Q_{x^{-1}} s^{-1}$ if x and s are invertible.
- (ii) $Q_{Q_s x} = Q_s Q_x Q_s$.

Lemma 3.4 ([16]) Let $x, s, p \in \Omega$. Then

- (i) $Q_{x^{1/2}}s$ and $Q_{s^{1/2}}x$ have the same eigenvalues.
- (ii) $Q_{x^{1/2}s}$ and $Q_{(O_px)^{1/2}}(Q_{p^{-1}s})$ have the same eigenvalues.

The following proposition was given by Vieira in [20] (see Proposition 5.6 in [20]), but we provide its proof using Lemma 3.3 and Lemma 3.4.

Proposition 3.1 Let Ψ be the proximity function defined in (13). Then we have

$$\Psi(\nu_{+}) = \Psi((Q_{(\nu+\alpha\,dx)^{1/2}}(\nu+\alpha\,ds))^{1/2})$$

Proof From $Q_{p^{-1}s} = Q_p x$ and (i) in Lemma 3.4, we know that $Q_p x$ and $Q_{x^{1/2}}p^2$ have the same eigenvalues. By the definition of p and (ii) in Lemma 3.3,

$$Q_{x^{1/2}}p^2 = Q_{x^{1/2}} \left(Q_x^{1/2} (Q_{x^{1/2}}s)^{-1/2} \right)^{-1} = Q_{x^{1/2}} Q_{x^{-1/2}} (Q_{x^{1/2}}s)^{1/2} = (Q_{x^{1/2}}s)^{1/2}$$

Then we can find $Q_{p_+^{-1}s_+}$ and $(Q_{x_+^{1/2}s_+})^{1/2}$ have the same eigenvalues. Here, $\sqrt{\mu}v_+ = Q_{p_+^{-1}s_+}$. We know that $x_+ = \sqrt{\mu}Q_{p^{-1}}(\nu + \alpha \, dx)$ and $s_+ = \sqrt{\mu}Q_p(\nu + \alpha \, ds)$, by the definition (8) and by (ii) in Lemma 3.4, then $(Q_{x_+^{1/2}s_+})^{1/2} = \sqrt{\mu}(Q(Q_{p^{-1}}(\nu + \alpha \, dx))^{1/2}(Q_p(\nu + \alpha \, ds)))^{1/2}$ and $\sqrt{\mu}(Q_{(\nu+\alpha \, dx)^{1/2}}(\nu + \alpha \, ds))^{1/2}$ have the same eigenvalues. Therefore, the proximity function satisfies the equality.

Then Proposition 2.1 and Proposition 3.1 imply the following inequality:

$$\Psi(v_+) \leq \frac{1}{2}\Psi(v+\alpha\,dx) + \frac{1}{2}\Psi(v+\alpha\,ds).$$

So, we can define $f_1(\alpha)$:

$$f(\alpha) \leq f_1(\alpha) := \frac{1}{2} \big(\Psi(\nu + \alpha \, dx) + \Psi(\nu + \alpha \, ds) \big) - \Psi(\nu).$$

To facilitate the forthcoming analysis, we also define, for any $x \in V$,

 $\lambda_{\min}(x) := \min \{ \lambda_i(x) \mid i = 1, \dots, r \}.$

The following lemma is obtained from Lemma 14 in [16] so that we can get the common lower bound of eigenvalues of $v + \alpha \, dx$ and $v + \alpha \, ds$, where α satisfies $v + \alpha \, dx \in \Omega$ and $v + \alpha \, ds \in \Omega$.

Lemma 3.5 For any $\alpha \in (0, \frac{\lambda_{\min}(\nu)}{\sigma})$,

 $\lambda_{\min}(v + \alpha \, dx) \ge \lambda_{\min}(v) - \alpha \sigma$ and $\lambda_{\min}(v + \alpha \, ds) \ge \lambda_{\min}(v) - \alpha \sigma$,

where σ is a number defined in (15).

Proof Let α be a fixed number in $(0, \frac{\lambda_{\min}(v)}{\sigma})$. From Lemma 14 in [16],

$$\lambda_{\min}(\nu + \alpha \, dx) \ge \lambda_{\min}(\nu) - \alpha \| dx \|.$$

Since $\sigma \ge \|dx\|$, we have

$$\lambda_{\min}(\nu + \alpha \, dx) \geq \lambda_{\min}(\nu) - \alpha \sigma.$$

Similarly, we obtain

$$\lambda_{\min}(\nu + \alpha \, ds) \ge \lambda_{\min}(\nu) - \alpha \sigma. \qquad \Box$$

The proof of the following proposition can be found in [18], but for the completeness, we give its detailed proof.

Proposition 3.2 ([18]) Suppose that the functions $\psi(x)$ and $\Psi(x)$ are defined by (2) and (13), respectively. Then, for any $\alpha \in (0, \frac{\lambda \min(v)}{\sigma})$,

$$\begin{aligned} \frac{d}{d\alpha}f_1(\alpha) &= \frac{1}{2}\operatorname{tr}\left(\psi'(\nu+\alpha\,dx)\circ\,dx\right) + \frac{1}{2}\operatorname{tr}\left(\psi'(\nu+\alpha\,ds)\circ ds\right),\\ \frac{d^2}{d\alpha^2}f_1(\alpha) &\leq \frac{3}{2}\max\left\{\Delta\psi'\left(\lambda_i(\nu+\alpha\,dx),\lambda_j(\nu+\alpha\,dx)\right) \mid i,j=1,\ldots,r\right\} \|dx\|^2 \\ &+ \frac{3}{2}\max\left\{\Delta\psi'\left(\lambda_i(\nu+\alpha\,ds),\lambda_j(\nu+\alpha\,ds)\right) \mid i,j=1,\ldots,r\right\} \|ds\|^2,\end{aligned}$$

where

$$\Delta \psi' \big(\lambda_i(\cdot), \lambda_j(\cdot) \big) = \begin{cases} \psi''(\lambda_i(\cdot)) & \text{if } \lambda_i(\cdot) = \lambda_j(\cdot), \\ \frac{\psi'(\lambda_i(\cdot)) - \psi'(\lambda_j(\cdot))}{\lambda_i(\cdot) - \lambda_j(\cdot)} & \text{if } \lambda_i(\cdot) \neq \lambda_j(\cdot). \end{cases}$$

Proof Using Lemma 3.1 in [21], we have

$$\frac{d}{d\alpha}\psi(\nu+\alpha\,dx) = \sum_{i=1}^{r} \Delta\psi(\lambda_{i}(\nu+\alpha\,dx),\lambda_{i}(\nu+\alpha\,dx))\langle c_{i}(\nu+\alpha\,dx),dx\rangle c_{i}(\nu+\alpha\,dx) + \sum_{1\leq j< l\leq r} 4\Delta\psi(\lambda_{j}(\nu+\alpha\,dx),\lambda_{l}(\nu+\alpha\,dx))) \times c_{j}(\nu+\alpha\,dx) \circ (c_{l}(\nu+\alpha\,dx)\circ dx).$$
(17)

Then we have

$$\frac{d}{d\alpha}\operatorname{tr}(\psi(\nu+\alpha\,dx)) = \frac{d}{d\alpha}\langle\psi(\nu+\alpha\,dx),e\rangle$$
$$= \left(\frac{d}{d\alpha}\psi(\nu+\alpha\,dx),e\right) = \operatorname{tr}\left(\frac{d}{d\alpha}\psi(\nu+\alpha\,dx)\right)$$

$$= \operatorname{tr}\left(\sum_{i=1}^{r} \psi'(\lambda_{i}(\nu + \alpha \, dx))\langle c_{i}(\nu + \alpha \, dx), dx\rangle c_{i}(\nu + \alpha \, dx)\right)$$

(by associativity of trace)
$$= \sum_{i=1}^{r} \psi'(\lambda_{i}(\nu + \alpha \, dx))\langle c_{i}(\nu + \alpha \, dx), dx\rangle \operatorname{tr}(c_{i}(\nu + \alpha \, dx))$$
$$= \left(\sum_{i=1}^{r} \psi'(\lambda_{i}(\nu + \alpha \, dx))c_{i}(\nu + \alpha \, dx), dx\right) \operatorname{tr}(c_{i}(\nu + \alpha \, dx)).$$

Then from Baes [17, 30] we know that $tr(c_i(\nu + \alpha dx)) = 1$, and hence, from the definition (3), we get

$$\frac{d}{d\alpha}\operatorname{tr}(\psi(\nu+\alpha\,dx))=\operatorname{tr}(\psi'(\nu+\alpha\,dx)\circ dx).$$

Thus, we have

$$\frac{d}{d\alpha}f_1(\alpha) = \frac{1}{2}\operatorname{tr}(\psi'(\nu + \alpha\,dx) \circ dx) + \frac{1}{2}\operatorname{tr}(\psi'(\nu + \alpha\,ds) \circ ds).$$

So, the first equality holds.

For the second inequality, we will use (17) by replacing ψ by ψ' .

$$\begin{split} \frac{d^2}{d\alpha^2} \operatorname{tr} \left(\psi \left(\nu + \alpha \, dx \right) \right) \\ &= \frac{d}{d\alpha} \operatorname{tr} \left(\psi' \left(\nu + \alpha \, dx \right) \circ dx \right) \\ &= \operatorname{tr} \left(\left(\frac{d}{d\alpha} \psi' \left(\nu + \alpha \, dx \right) \right) \circ dx \right) \\ &= \operatorname{tr} \left(\left(\sum_{i=1}^r \Delta \psi' \left(\lambda_i (\nu + \alpha \, dx), \lambda_i (\nu + \alpha \, dx) \right) \left\langle c_i (\nu + \alpha \, dx), dx \right\rangle c_i (\nu + \alpha \, dx) \right) \\ &+ \sum_{1 \leq j < l \leq r} 4 \Delta \psi' \left(\lambda_j (\nu + \alpha \, dx), \lambda_l (\nu + \alpha \, dx) \right) \\ &\times c_j (\nu + \alpha \, dx) \circ \left(c_l (\nu + \alpha \, dx) \circ dx \right) \right) \circ dx \right). \end{split}$$

Here, let $dx = \sum_{j=1}^{r} \lambda_j(dx) c_j(dx)$. Then we have

$$\sum_{i=1}^{r} \left(\operatorname{tr} \left(c_i(\nu + \alpha \, dx) \circ dx \right) \right)^2$$
$$= \sum_{i=1}^{r} \left(\operatorname{tr} \left(\sum_{j=1}^{r} \lambda_j(dx) c_i(\nu + \alpha \, dx) \circ c_j(dx) \right) \right)^2$$
$$= \sum_{i=1}^{r} \left(\sum_{j=1}^{r} \lambda_j(dx) \operatorname{tr} \left(c_i(\nu + \alpha \, dx) \circ c_j(dx) \right) \right)^2.$$

Since $c_i(v + \alpha \, dx)$ and $c_i(dx)$ are in $\overline{\Omega}$ which is a self-dual cone, then

$$\operatorname{tr}(c_i(\nu+\alpha\,dx)\circ c_j(dx))\geq 0, \quad \text{for } i,j=1,\ldots,r.$$

Furthermore, $\sum_{j=1}^{r} \operatorname{tr}(c_i(\nu + \alpha \, dx) \circ c_j(dx)) = \operatorname{tr}(c_i(\nu + \alpha \, dx)) = 1$. Then we have

$$\sum_{i=1}^{r} \left(\operatorname{tr} \left(c_i(\nu + \alpha \, dx) \circ dx \right) \right)^2 = \sum_{i=1}^{r} \left(\sum_{j=1}^{r} \operatorname{tr} \left(c_i(\nu + \alpha \, dx) \circ c_j(dx) \right) \lambda_j(dx) \right)^2$$
$$\leq \sum_{i=1}^{r} \left(\sum_{j=1}^{r} \lambda_j^2(dx) \operatorname{tr} \left(c_i(\nu + \alpha \, dx) \circ c_j(dx) \right) \right)$$
$$= \sum_{j=1}^{r} \left(\lambda_j^2(dx) \sum_{i=1}^{r} \operatorname{tr} \left(c_i(\nu + \alpha \, dx) \circ c_j(dx) \right) \right).$$

Since $\sum_{i=1}^{r} \operatorname{tr}(c_i(\nu + \alpha \, dx) \circ c_j(dx)) = \operatorname{tr}(c_j(dx)) = 1$, we have

$$\sum_{i=1}^{r} \left(\operatorname{tr} \left(c_i (\nu + \alpha \, dx) \circ dx \right) \right)^2 \leq \| dx \|^2.$$
(18)

Now, we decompose dx along Lemma 1.1 such as $dx = \sum_{1 \le i \le k \le r} dx_{ik}$ for the system of idempotent $\{c_1(v + \alpha \, dx), \dots, c_r(v + \alpha \, dx)\}$. Then, for j < l,

$$\operatorname{tr}\left(\left(c_{j}(\nu+\alpha\,dx)\circ dx\right)\circ\left(c_{l}(\nu+\alpha\,dx)\circ dx\right)\right)$$

$$=\operatorname{tr}\left(\left(c_{j}(\nu+\alpha\,dx)\circ\sum_{1\leq i\leq k\leq r}dx_{ik}\right)\circ\left(c_{l}(\nu+\alpha\,dx)\circ\sum_{1\leq i\leq k\leq r}dx_{ik}\right)\right)$$

$$=\operatorname{tr}\left(\left(c_{j}(\nu+\alpha\,dx)\circ\sum_{i=1}^{r}dx_{ji}\right)\circ\left(c_{l}(\nu+\alpha\,dx)\circ\sum_{i=1}^{r}dx_{li}\right)\right)$$
(by (ii) in Lemma 1.1)
$$=\operatorname{tr}\left(\frac{1}{4}\,dx_{jl}\circ dx_{lj}\right)$$
(by (iii) and (iv) in Lemma 1.1)
$$=\frac{1}{4}\operatorname{tr}\left(dx_{jl}^{2}\right).$$

This means, for each j < l,

$$\operatorname{tr}((c_j(\nu+\alpha\,dx)\circ dx)\circ(c_l(\nu+\alpha\,dx)\circ dx))\geq 0.$$

Moreover, we have,

$$\sum_{1 \leq j < l \leq r} \operatorname{tr}\left(\left(c_{j}(\nu + \alpha \, dx) \circ dx\right) \circ \left(c_{l}(\nu + \alpha \, dx) \circ dx\right)\right)$$
$$\leq \sum_{1 \leq j < l \leq r} \operatorname{tr}\left(\left(c_{j}(\nu + \alpha \, dx) \circ dx\right) \circ \left(c_{l}(\nu + \alpha \, dx) \circ dx\right)\right) + \frac{1}{2} \sum_{i=1}^{r} \operatorname{tr}\left(\left(c_{i}(\nu + \alpha \, dx) \circ dx\right)^{2}\right)$$

$$= \frac{1}{2} \operatorname{tr} \left(\left(\sum_{i=1}^{r} \left(c_i (\nu + \alpha \, dx) \circ dx \right) \right)^2 \right)$$
$$= \frac{1}{2} \operatorname{tr} \left(dx^2 \right)$$
$$= \frac{1}{2} \| dx \|^2.$$
(19)

Since for each *i*, $(tr(c_i(\nu + \alpha dx) \circ dx))^2$ are nonnegative and for each *j*, *l* with *j* < *l*, $tr((c_j(\nu + \alpha dx) \circ dx) \circ (c_l(\nu + \alpha dx) \circ dx))$ are nonnegative, we get from (18) and (19)

$$\frac{d^2}{d\alpha^2}\operatorname{tr}(\psi(\nu+\alpha\,dx)) \leq 3\max\left\{\Delta\psi'(\lambda_i(\nu+\alpha\,dx),\lambda_j(\nu+\alpha\,dx)) \mid i,j=1,\ldots,r\right\} \|dx\|^2.$$

Similarly,

$$\frac{d^2}{d\alpha^2}\operatorname{tr}(\psi(\nu+\alpha\,ds)) \leq 3\max\{\Delta\psi'(\lambda_i(\nu+\alpha\,ds),\lambda_j(\nu+\alpha\,ds)) \mid i,j=1,\ldots,r\}\|ds\|^2.$$

From the definition of $f_1(\alpha)$,

$$\frac{d^2}{d\alpha^2}f_1(\alpha) = \frac{d^2}{d\alpha^2}\left(\frac{1}{2}\operatorname{tr}(\psi(\nu+\alpha\,dx)) + \frac{1}{2}\operatorname{tr}(\psi(\nu+\alpha\,ds))\right).$$

Thus, we have the conclusion.

The next result presents an upper bound for the second derivative of $f_1(\alpha)$ which is usable for establishing the polynomial complexity of the algorithm.

Proposition 3.3 *For any* $\alpha \in (0, \frac{\lambda_{\min}(\nu)}{\sigma})$ *,*

$$f_1''(\alpha) \leq \frac{3}{2} \psi'' (\lambda_{\min}(\nu) - \alpha \sigma) \sigma^2.$$

Proof Since $\psi''(t)$ is a decreasing function on $t \in (0, \infty)$, using Lemma 3.5 and the mean value theorem, we have

$$\psi''(\lambda_{\min}(\nu) - \alpha \sigma) \ge \max\{\Delta \psi'(\lambda_i(\nu + \alpha \, dx), \lambda_j(\nu + \alpha \, dx)) \mid i, j = 1, \dots, r\}$$

and

$$\psi''(\lambda_{\min}(\nu) - \alpha \sigma) \geq \max \{ \Delta \psi'(\lambda_i(\nu + \alpha \, ds), \lambda_j(\nu + \alpha \, ds)) \mid i, j = 1, \dots, r \}.$$

Thus, by Proposition 3.2,

$$\begin{aligned} \frac{d^2}{d\alpha^2} f_1(\alpha) &\leq \frac{3}{2} \psi'' \big(\lambda_{\min}(\nu) - \alpha \sigma \big) \| dx \|^2 + \frac{3}{2} \psi'' \big(\lambda_{\min}(\nu) - \alpha \sigma \big) \| ds \|^2 \\ &= \frac{3}{2} \psi'' \big(\lambda_{\min}(\nu) - \alpha \sigma \big) \sigma^2. \end{aligned}$$

We can easily check that $f_1(0) = 0$ and $f'_1(0) = -\frac{\sigma^2}{2}$. By Proposition 3.3, we obtain an upper bound $f_2(\alpha)$ for $f_1(\alpha)$ as follows:

$$f_{1}(\alpha) = f_{1}(0) + f_{1}'(0)\alpha + \int_{0}^{\alpha} \int_{0}^{\xi} f_{1}''(\zeta) d\zeta d\xi$$
$$\leq f_{2}(\alpha) := f_{1}(0) + f_{1}'(0)\alpha + \frac{3}{2}\sigma^{2} \int_{0}^{\alpha} \int_{0}^{\xi} \psi''(\lambda_{\min}(\nu) - \alpha\sigma) d\zeta d\xi.$$

Note that $f_2(0) = 0$. Furthermore, since $f'_2(\alpha) = -\frac{\sigma^2}{2} + \frac{3\sigma}{2}(\psi'(\lambda_{\min}(\nu)) - \psi'(\lambda_{\min}(\nu) - \alpha\sigma))$, we have $f'_2(0) = -\frac{\sigma^2}{2}$ which is the same value of $f'_1(0)$, and $f''_2(\alpha) = \frac{3\sigma^2}{2}\psi''(\lambda_{\min}(\nu) - \alpha\sigma)$ which is increasing on $\alpha \in [0, \frac{\lambda_{\min}(\nu)}{\sigma})$. Using $f'_1(0) = f'_2(0)$ and $f''_1(\alpha) \leq f''_2(\alpha)$, we can easily check that

$$f_1'(\alpha) = f_1'(0) + \int_0^{\alpha} f_1''(\xi) d\xi \leq f_2'(\alpha).$$

This relation gives that

$$f_1'(\alpha) \leq 0$$
, if $f_2'(\alpha) \leq 0$.

To compute the feasible step size α such that the proximity measure is decreasing when we take a new iterate for fixed μ , we want to calculate the step size α which satisfies that $f'_2(\alpha) \leq 0$ holds with α as large as possible. Since $f''_2(\alpha) > 0$, that is, $f'_2(\alpha)$ is monotonically increasing at α , the largest possible value at α satisfying $f'_2(\alpha) \leq 0$ occurs when $f'_2(\alpha) = 0$, that is,

$$-\psi'(\lambda_{\min}(\nu) - \alpha\sigma) + \psi'(\lambda_{\min}(\nu)) = \frac{\sigma}{3}.$$
 (20)

Since $\psi''(t)$ is monotonically decreasing, the derivative of the left-hand side in (20) with respect to $\lambda_{\min}(v)$ is

$$-\psi''(\lambda_{\min}(\nu)-\alpha\sigma)+\psi''(\lambda_{\min}(\nu))<0.$$

So, the left-hand side in (20) is decreasing at $\lambda_{\min}(\nu)$. This implies that if $\lambda_{\min}(\nu)$ becomes smaller, then α gets smaller with fixed σ . Note that

$$\sigma = \sqrt{\sum_{i=1}^{n} (\psi'(\lambda_i(\nu)))^2} \ge |\psi'(\lambda_{\min}(\nu))| \ge -\psi'(\lambda_{\min}(\nu))$$

and the equality is true if and only if $\lambda_{\min}(\nu)$ is the only coordinate in $(\lambda_1(\nu), \dots, \lambda_r(\nu))$ which is different from 1 and $\lambda_{\min}(\nu) < 1$, that is, $\psi'(\lambda_{\min}(\nu)) < 0$. Hence, the worse situation for the largest step size occurs when $\lambda_{\min}(\nu)$ satisfies

$$-\psi'(\lambda_{\min}(\nu)) = \sigma.$$
⁽²¹⁾

In that case, the largest α satisfying (20) is minimal. For our purpose, we need to deal with the worse case, and so we assume that (21) holds.

From now on, we denote that $\rho : [0, \infty) \to (0, 1]$ is the inverse function of the restriction of $-\psi'(t)$ in the interval (0, 1]. Then (21) implies

$$\lambda_{\min}(\nu) = \rho(\sigma). \tag{22}$$

By using (20) and (21), we immediately obtain

$$-\psi'(\lambda_{\min}(\nu)-\alpha\sigma)=\frac{4}{3}\sigma.$$

By the definition of ρ and (22), the largest step size α of the worse case is given as follows:

$$\alpha^* = \frac{\rho(\sigma) - \rho(\frac{4}{3}\sigma)}{\sigma}.$$
(23)

For the purpose of finding an upper bound of $f(\alpha)$, we need a default step size $\bar{\alpha}$ that is the lower bound of the α^* and consists of σ .

Lemma 3.6 Let $\sigma \ge 1$. Then, for $0 < t \le \rho(\frac{4}{3}\sigma)$,

$$\psi''(t) \leq 1 + 3\sigma(1 + pq + q) \left(1 + \frac{1}{p}\log 3\sigma\right)^{\frac{q+1}{q}}.$$

Proof From $\psi'(t) = t - t^{-q-1} \cdot e^{p(t^{-q}-1)}$, let $-\psi'_b(t) = t^{-q-1} \cdot e^{p(t^{-q}-1)}$ and let $\underline{\rho} : [1, \infty) \to (0, 1]$ denote the inverse function of the restriction of $-\psi'_b(t)$ to the interval (0, 1]. Let $\rho(\frac{4}{3}\sigma) = \tilde{t}$. Then $0 < \tilde{t} \leq 1$ and $\frac{4}{3}\sigma = -\psi'(\tilde{t}) = -\tilde{t} - \psi'_b(\tilde{t})$. So, $-\psi'_b(\tilde{t}) = \tilde{t} + \frac{4}{3}\sigma \leq 1 + 2\sigma \leq 3\sigma$. Since $\underline{\rho}$ is a decreasing function, $(\rho(\frac{4}{3}\sigma) = \tilde{t} =)\underline{\rho}(-\psi'_b(\tilde{t})) \geq \underline{\rho}(3\sigma)$. Let $\underline{\rho}(3\sigma) = \hat{t}$. Then

$$3\sigma = -\psi_b'(\hat{t}) = (\underline{\rho}(3\sigma))^{-q-1} \cdot e^{p((\underline{\rho}(3\sigma))^{-q}-1)}$$
(24)

implies

$$e^{p((\underline{\rho}(3\sigma))^{-q}-1)} = 3\sigma(\underline{\rho}(3\sigma))^{q+1} \leq 3\sigma \quad \Rightarrow \quad p((\underline{\rho}(3\sigma))^{-q}-1) \leq \log 3\sigma$$

$$\Rightarrow \quad \underline{\rho}(3\sigma) \geq \left(1 + \frac{1}{p}\log 3\sigma\right)^{-\frac{1}{q}}, \tag{25}$$

$$\psi''(\hat{t}) = 1 + ((q+1)\hat{t}^{q} + pq)\hat{t}^{-q-1} \cdot e^{p(\hat{t}^{-q}-1)} \cdot \hat{t}^{-q-1} = 1 + ((q+1)\hat{t}^{q} + pq)(-\psi_{b}'(\hat{t})) \cdot \hat{t}^{-q-1}$$

$$= 1 + ((q+1)\hat{t}^{q} + pq) \cdot 3\sigma \cdot (\underline{\rho}(3\sigma))^{-q-1} \leq 1 + 3\sigma(1 + pq + q)\left(1 + \frac{1}{p}\log 3\sigma\right)^{\frac{q+1}{q}},$$

the last inequality comes from $\hat{t} \in (0, 1]$ and (25).

Now, we present a lower bound of the value of α^* .

Theorem 3.2 Let α^* be as defined in (23). Then

$$\alpha^* \ge \frac{1}{3(1+3\sigma(1+pq+q)(1+\frac{1}{p}\log 3\sigma)^{\frac{q+1}{q}})}.$$

Proof Since $-\psi'(\rho(\sigma)) = \sigma$, taking the derivative of σ at both sides, we get

$$\rho'(\sigma) = -\frac{1}{\psi''(\rho(\sigma))}.$$

Moreover, we have

$$\alpha^* = \frac{1}{\sigma} \int_{\frac{4}{3}\sigma}^{\sigma} \rho'(\xi) d\xi = \frac{1}{\sigma} \int_{\sigma}^{\frac{4}{3}\sigma} \frac{1}{\psi''(\rho(\xi))} d\xi \ge \frac{1}{\sigma} \left[\frac{\xi}{\psi''(\rho(\frac{4}{3}\sigma))} \right]_{\sigma}^{\frac{4}{3}\sigma} = \frac{1}{3\psi''(\rho(\frac{4}{3}\sigma))},$$

where the inequality follows from $\sigma \leq \xi \leq \frac{4}{3}\sigma$ and ρ and ψ'' are monotonically decreasing. Also, by Lemma 3.6, we can complete the proof.

For using $\bar{\alpha}$ as the default step size in the algorithm for the SOP, define the $\bar{\alpha}$ as follows:

$$\bar{\alpha} = \frac{1}{3(1+3\sigma(1+pq+q)(1+\frac{1}{p}\log 3\sigma)^{\frac{q+1}{q}})}.$$
(26)

We will use $\bar{\alpha}$ as the default step size in our algorithm.

3.3 Decrease of the proximity function during an inner iteration

Now, we show that our proximity function Ψ with our default step size $\bar{\alpha}$ is decreasing. It can be easily established by using the following result.

Lemma 3.7 ([4]) Let h(t) be a twice differentiable convex function with h(0) = 0, h'(0) < 0and let h(t) attain its (global) minimum at $t^{\circ} > 0$. If h''(t) is increasing for $t \in [0, t^{\circ}]$, then

$$h(t) \leq rac{th'(0)}{2}, \quad 0 \leq t \leq t^*.$$

Since $f_2(\alpha)$ satisfies assumptions of the above lemma,

$$f(\alpha) \leq f_1(\alpha) \leq f_2(\alpha) \leq \frac{f_2'(0)}{2} \alpha \quad \text{for all } 0 \leq \alpha \leq \alpha^*.$$

Since $f'_2(0) = -\frac{\sigma^2}{2}$, we can obtain the upper bound for the decreasing value of the proximity in the inner iteration by Lemma 3.7.

Theorem 3.3 Let $\bar{\alpha}$ be the default step size as defined in (26). Then we have

$$f(\bar{\alpha}) \leq -\frac{1}{6} \cdot \frac{\sqrt{\Psi}}{1 + 3\sqrt{2}(1 + pq + q)(1 + \frac{1}{p}\log 3\sqrt{2\Psi_0})^{\frac{q+1}{q}}}$$

Proof Since $f'_2(0) = -\frac{\sigma^2}{2}$ and $\bar{\alpha} \in [0, \alpha^*]$, we have

$$f(\bar{\alpha}) \leq \frac{1}{2}\bar{\alpha}f_2'(0) = \frac{1}{2} \cdot \frac{1}{3(1+3\sigma(1+pq+q)(1+\frac{1}{p}\log 3\sigma)^{\frac{q+1}{q}})} \cdot \left(-\frac{\sigma^2}{2}\right)$$
$$= -\frac{1}{12} \cdot \frac{\sigma^2}{1+3\sigma(1+pq+q)(1+\frac{1}{p}\log 3\sigma)^{\frac{q+1}{q}}}.$$

This expresses the decrease in one inner iteration in terms of σ . Since the decrease depends monotonically on σ , we can express the decrease in terms of $\Psi = \Psi(\nu)$ by Lemma 2.1 as follows:

$$f(\bar{\alpha}) \leq -\frac{1}{6} \cdot \frac{\Psi}{1 + 3\sqrt{2\Psi}(1 + pq + q)(1 + \frac{1}{p}\log 3\sqrt{2\Psi})^{\frac{q+1}{q}}} \\ \leq -\frac{1}{6} \cdot \frac{\sqrt{\Psi} \cdot \sqrt{\Psi}}{\sqrt{\Psi} + 3\sqrt{2\Psi}(1 + pq + q)(1 + \frac{1}{p}\log 3\sqrt{2\Psi_0})^{\frac{q+1}{q}}} \\ = -\frac{1}{6} \cdot \frac{\sqrt{\Psi}}{1 + 3\sqrt{2}(1 + pq + q)(1 + \frac{1}{p}\log 3\sqrt{2\Psi_0})^{\frac{q+1}{q}}},$$

where the inequality follows from $\Psi_0 \geqq \Psi \geqq \tau \geqq 1$. The theorem is satisfied.

3.4 Iteration bound

We need to count how many inner iterations are required to return to the situation where $\Psi(\nu) \leq \tau$ after a μ -update. We denote the value of $\Psi(\nu)$ after μ -update as Ψ_0 ; the subsequent values in the same outer iteration are denoted as Ψ_k , k = 1, ... If K denotes the total number of inner iterations in the outer iteration, then we have

$$\Psi_0 \leq L(r, \theta, \tau) = \mathcal{O}(r), \qquad \Psi_{K-1} > \tau, \qquad 0 \leq \Psi_K \leq \tau$$

and according to Theorem 3.3,

$$\Psi_{k+1} \leq \Psi_k - \frac{1}{6 + 18\sqrt{2}(1 + pq + q)(1 + \frac{1}{p}\log 3\sqrt{2\Psi_0})^{\frac{q+1}{q}}} \Psi_k^{\frac{1}{2}}.$$

At this stage, we invoke Lemma 14 in [4].

Lemma 3.8 ([4]) Let t_0, t_1, \ldots, t_K be a sequence of positive numbers such that

$$t_{k+1} \leq t_k - \beta t_k^{1-\gamma}, \quad k = 0, 1, \dots, K-1,$$

where $\beta > 0$ and $0 < \gamma \leq 1$. Then

$$K \leq rac{t_0^{\gamma}}{\beta \gamma}.$$

Letting $t_k = \Psi_k$, $\beta = \frac{1}{\frac{6+18\sqrt{2}(1+pq+q)(1+\frac{1}{p}\log 3\sqrt{2\Psi_0})}{q+1}}$ and $\gamma = \frac{1}{2}$, we can get the following lemma from Lemma 3.8.

Lemma 3.9 *Let K be the total number of inner iterations in the outer iteration. Then we have*

$$K \leq 2\left(6 + 18\sqrt{2}(1 + pq + q)\left(1 + \frac{1}{p}\log 3\sqrt{2}\sqrt{\Psi_0}\right)^{\frac{q+1}{q}}\right)\Psi_0^{1/2},$$

where Ψ_0 is the value of $\Psi(v)$ after the μ -update in the outer iteration.

Now, we estimate the total number of iterations of our algorithm.

Theorem 3.4 If $\tau \ge 1$ and $0 < \theta < 1$, the total number of iterations is not more than

$$\left\lceil 2\left(6+18\sqrt{2}(1+pq+q)\left(1+\frac{1}{p}\log 3\sqrt{2}\sqrt{\Psi_0}\right)^{\frac{q+1}{q}}\right)\Psi_0^{1/2}\right\rceil \left\lceil \frac{1}{\theta}\log \frac{r}{\epsilon}\right\rceil.$$

Proof In the algorithm, $r\mu \ge \epsilon$, $\mu_k := (1 - \theta)^k \mu_0$ and $\mu_0 = 1$. By simple computation, we have

$$k \leq \frac{1}{\theta} \log \frac{r}{\epsilon}.$$

Therefore, the number of outer iterations is bounded above by

$$\frac{1}{\theta}\log\frac{r}{\epsilon}.$$

Multiplication of this result by the number in the above lemma satisfies the theorem. $\hfill\square$

Since $\Psi_0^{1/2} = \mathcal{O}(\sqrt{r})$, if we take $p = \mathcal{O}(\log r)$ and q = 1, then we can get the best known upper bound for the total number of inner iterations in the outer iteration is

 $\mathcal{O}(\sqrt{r}\log r).$

Also, we take for θ a constant (not depending on r), namely $\frac{1}{\theta} = \Theta(1)$. With $\tau = O(r)$, the best complexity of the primal-dual interior-point method for a linear optimization problem based on our new proximity function with $p = \log r$ and q = 1 is given by

$$\mathcal{O}\left(\sqrt{r}\log r\log\frac{r}{\epsilon}\right)$$

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

The authors, together discussed and solved the problems in the manuscript. All authors read and approved the final manuscript.

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